

Scattering of electromagnetic wave.  
① Scattering by free charge or Thomson's scattering.  
Let us consider an electron of mass  $m$  and charge  $e$  in the path of polarised light.

Let the incident E.M wave having electric intensity  $E$  along  $x$ -direction moving along  $z$ -direction. The electromagnetic force, is  $e \cdot E$ .

The equation of motion of electron is given by

$$m \cdot \frac{d^2x}{dt^2} = eE = eE_0 e^{-i(\omega t - kx)} \quad \text{--- (A)}$$

$$\text{or, } \frac{d^2x}{dt^2} = \frac{eE_0}{m} e^{-i(\omega t - kx)} \quad \text{--- (B)}$$

$$\Rightarrow x = -\frac{eE_0}{\omega^2 m} e^{-i(\omega t - kx)}$$

The oscillating charge behaves like a dipole having dipole moment

$$p = ex = -\frac{e^2 E_0}{\omega^2 m} e^{-i(\omega t - kx)}$$

$$p = p_0 e^{-i(\omega t - kx)}$$

$$\text{where } p_0 = -\frac{e^2 E_0}{m \omega^2}$$

The oscillating dipole radiates energy. The average energy radiating per second per unit area is given by

$$\omega_s = \frac{1}{4\pi\epsilon_0} \frac{\omega^4 p_0^2}{8\pi c^3 r^2} \sin^2 \alpha$$

$$= \frac{1}{4\pi\epsilon_0} \frac{e^4 E_0^2}{8\pi c^3 r^2 m^2} \sin^2 \alpha \quad \text{--- (C)}$$

For plane electromagnetic wave the average incident energy is given by

$$\omega_i = \frac{1}{2} c \epsilon_0 E_0^2$$

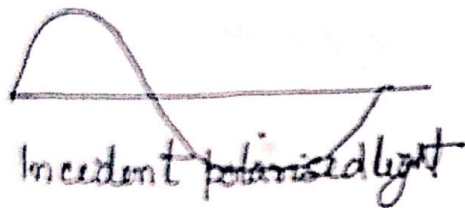
The differential scattering cross-section is given by

$$\begin{aligned} \sigma(\theta) &= \frac{\omega_s}{\omega_i} r^2 \\ &= \frac{1}{4\pi\epsilon_0} \frac{e^4 E_0^2 \sin^2 \alpha}{8\pi c^3 r^2 m^2} \cdot \frac{2}{c \epsilon_0 E_0^2} r^2 \end{aligned}$$

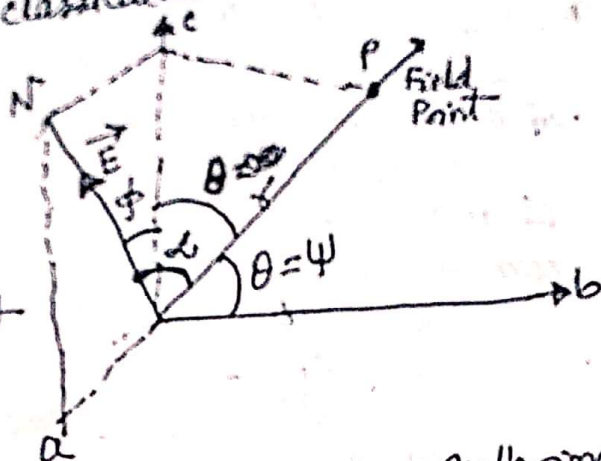
$$\text{or } \sigma(\theta) = \left( \frac{e^2}{4\pi\epsilon_0 mc^2} \right)^2 \sin^2 \alpha \quad (2)$$

$$\text{or } \sigma(\theta) = r_0^2 \sin^2 \alpha$$

where  $r_0 = \frac{e^2}{4\pi\epsilon_0 mc^2} =$  classical radius of the electron.



Incident polarised light



Let  $\phi$  is the polarizing angle and  $\theta$  is the scattering angle of e.m. wave and  $\alpha$  is the angle between position vector  $r$  & electric vector  $E$

$$\text{Now } ON = r \cos \alpha \quad \& \quad OM = r \sin \theta$$

$$\therefore ON = r \cos \alpha = OM \cos \phi = r \sin \theta \cos \phi$$

$$\text{Again } \cos^2 \alpha = \sin^2 \theta \cdot \cos^2 \phi$$

$$\therefore 1 - \sin^2 \alpha = (1 - \cos^2 \theta) \cos^2 \phi$$

$$\text{or } \sin^2 \alpha = 1 - \cos^2 \phi (1 - \cos^2 \theta)$$

For plane polarised light,  $\phi = 0$

$$\therefore \sin^2 \alpha = \cos^2 \theta$$

For unpolarised light, the average value of  $\phi$  is equal to  $\sqrt{1/2}$

$$\therefore \sin^2 \alpha = 1 - \frac{1}{2} (1 - \cos^2 \theta) = \frac{1}{2} (1 + \cos^2 \theta)$$

$$= \frac{1}{2} (1 + \cos^2 \theta)$$

The scattering cross-section area is given by

$$\boxed{\sigma(\theta) = r_0^2 \frac{1}{2} (1 + \cos^2 \theta)}$$

This is Thomson scattering formula.

(c) Scattering by bounded charge (Rayleigh scattering)  
 Let us consider following forces are acting upon the bounded electrons.

(i) Inertial force =  $m \frac{d^2x}{dt^2}$

(ii) Frictional force which is proportional to velocity  
 $f_r = \gamma \frac{dx}{dt}$

(iii) Damping force which is proportional to the displacement  $f_s = Kx$

(iv) The periodic electromagnetic force.  
 $F = eE_0 e^{i(\omega t - Kx)}$

The eqn of motion is given by.

$$m \frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + Kx = eE_0 e^{-i(\omega t - Kx)}$$

$$\text{or, } \frac{d^2x}{dt^2} + \frac{\gamma}{m} \frac{dx}{dt} + \frac{K}{m} x = \frac{e}{m} E_0 e^{-i(\omega t - Kx)} \quad \text{--- (A)}$$

$$\text{or, } \frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = \frac{e}{m} E_0 e^{-i(\omega t - Kx)}$$

where  $\gamma = \frac{\gamma}{m} \Rightarrow \omega_0^2 = \frac{K}{m}$ .

The solution of complementary function is given by. --- (B)

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = 0$$

Let  $x = A e^{\alpha t}$ , the eqn's may be written as

$$\alpha^2 + \gamma \alpha + \omega_0^2 = 0$$

$$\therefore \alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - \omega_0^2}$$

$$= -\frac{\gamma}{2} \pm i \sqrt{\omega_0^2 - \frac{\gamma^2}{4}}$$

$$\therefore x = e^{-\gamma t/2} [A_1 e^{i(\omega_0^2 - \gamma^2/4)t} + A_2 e^{-i(\omega_0^2 - \gamma^2/4)t}] \quad \text{--- (C)}$$

The factor decreases exponentially which die after some time. The sol<sup>n</sup> of perpendicular integral is given by

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = \frac{e}{m} E_0 e^{-i(\omega t - Kx)} \quad \text{--- (D)}$$

Let  $x = \beta e^{-i(\omega t - Kx)}$  (4)

From eqn (1) we get

$$\beta = \frac{eE_0}{m(\omega_0^2 - \omega^2 - i\gamma\omega)} e^{-i(\omega t - Kx)}$$

$$\begin{aligned} x &= \frac{eE_0}{m(\omega_0^2 - \omega^2 - i\gamma\omega)} e^{-i(\omega t - Kx)} \\ &= \frac{eE_0 [(\omega_0^2 - \omega^2) + i\gamma\omega]}{m[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]} e^{-i(\omega t - Kx)} \\ &= \frac{eE_0}{m} \frac{e^{-i(\omega t - Kx)}}{[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} \left[ \frac{\omega_0^2 - \omega^2}{[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} + \frac{i\gamma\omega}{[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} \right] \\ &= \frac{eE_0}{m} \frac{e^{-i(\omega t - Kx)}}{[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} [\cos\theta - i\sin\theta] \\ &= \frac{eE_0}{m} \frac{e^{-i(\omega t - Kx) + \delta}}{[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} \quad \text{--- (E)} \end{aligned}$$

where  $\delta = \tan^{-1} \frac{\gamma\omega}{\omega_0^2 - \omega^2}$

The complete solution of equation (A) is given by

$$x = e^{-\gamma t/2} [A_1 e^{i\beta t} + A_2 e^{-i\beta t}] + \frac{eE_0}{m[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} e^{i(\omega t - Kx + \delta)} \quad \text{--- (F)}$$

The first term may be neglected due to -ve exponential  
the general eqn of oscillating charge due to dipole moment

$$p = e x = \frac{e^2 E_0}{m[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} [e^{i(\omega t - Kx + \delta)}]$$

or,  $p = p_0 e^{-i(\omega t - Kx - \delta)} \quad \text{--- (G)}$

where  $p_0 = \frac{e^2 E}{m[(\omega_0^2 - \omega^2)^2 + \gamma^2\omega^2]^{1/2}} \quad \text{--- (H)}$

The average energy radiated per second per unit area is given by.

$$S = \frac{1}{4\pi\epsilon_0} \frac{\omega^4}{8\pi c^2 r^2} \cdot p_0^2 \sin^2 \theta \quad \text{--- (ii)}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{\omega^4 p^2}{8\pi c^2 r^2} \frac{e^4 E_0^2 \sin^2 \theta}{m[(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2]} \quad \text{--- (ii)}$$

The average incident radiation (Poynting Vector)

$$S_{in} = \frac{1}{2} E_0 C E_0^2$$

The differential scattering cross section is given by

$$\sigma(\theta) = \frac{S}{S_{in}} \omega^2 \gamma^2$$

$$= \frac{1}{4\pi\epsilon_0} \frac{\omega^4}{8\pi c^2 r^2} \frac{e^2 E_0^2 \sin^2 \theta \gamma^2}{m[(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2] C E_0 E_0^2}$$

on solving we get

$$\sigma(\theta) = \frac{2}{3} \gamma_0^2 \omega^4 \left( \frac{1 + \cos^2 \theta}{2} \right) \frac{1}{(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2} \quad \text{--- (I)}$$

The total scattering cross section is given by

$$\sigma_T = \int_0^\pi \frac{2}{3} \gamma_0^2 \omega^4 \left( \frac{1 + \cos^2 \theta}{2} \right) \cdot 2\pi \sin \theta d\theta$$

$$\sigma_T = \frac{8\pi}{3} \frac{\gamma_0^2 \omega^4}{[(\omega_0^2 - \omega^2)^2 + \gamma^2 \omega^2]}$$

The eqn I represents the total scattering cross section for elastically bound electron. The scattering cross section is a function of frequency of incident radiation is given by Graph A.

